

On some properties of the Unitary Dual Group

Michaël Ulrich

Université de Franche-Comté (France)

and

Ernst-Moritz-Arndt Universität Greifswald (Germany)

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Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$

A stochastic process on
 $U(n)$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Definition (Dual group)

A (unital) dual group $(\mathcal{A}, \Delta, \delta, \Sigma)$ is a (unital) $*$ -algebra \mathcal{A} , and three (unital) $*$ -homomorphisms $\Delta : \mathcal{A} \rightarrow \mathcal{A} \sqcup \mathcal{A}$, $\delta : \mathcal{A} \rightarrow \mathbb{C}$ and $\Sigma : \mathcal{A} \rightarrow \mathcal{A}$, such that

- ▶ The map Δ is a coassociative coproduct:

$$(Id \sqcup \Delta) \circ \Delta = (\Delta \sqcup Id) \circ \Delta$$
- ▶ The map δ is a counit: $(\delta \sqcup Id) \circ \Delta = Id = (Id \sqcup \delta) \circ \Delta$
- ▶ The map Σ is a coinverse:

$$(\Sigma \sqcup Id) \circ \Delta = 1_{\mathcal{A}} \circ \delta = (Id \sqcup \Sigma) \circ \Delta.$$

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U(n)$ A stochastic process on $U(n)$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

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$$(\Sigma \sqcup Id) \circ \Delta = 1_{\mathcal{A}} \circ \delta = (Id \sqcup \Sigma) \circ \Delta.$$

Remark

- ▶ Free product "more noncommutative" than tensor product
- ▶ Unlike compact quantum groups, we have a counit and a coinverse (cf. compact quantum groups of Kac type)

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

This is the dual group we will be interested in the sequel

Unitary Dual
Group

Michaël Ulrich

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$

A stochastic process on
 $U(n)$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Definition

Let $n \geq 0$. We call unitary dual group the following dual group: $(U\langle n \rangle, \Delta, \delta, \Sigma)$ where:

- ▶ The unital $*$ -algebra $U\langle n \rangle$ is generated by n^2 elements $(u_{ij})_{1 \leq i, j \leq n}$ with the relations

$$\sum_k u_{ki}^* u_{kj} = \delta_{ij} = \sum_k u_{ik} u_{jk}^*.$$

- ▶ The coproduct is given on the generators by $\Delta(u_{ij}) = \sum_k u_{ik}^{(1)} u_{kj}^{(2)}$.
- ▶ The counit is given by $\delta(u_{ij}) = \delta_{ij}$.
- ▶ The antipode is given by: $\Sigma(u_{ij}) = u_{ji}^*$.

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

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- ▶ The antipode is given by: $\Sigma(u_{ij}) = u_{ji}^*$.

The relations defining $U\langle n \rangle$ can be summed up by saying that $u = (u_{ij})_{1 \leq i, j \leq n}$ is a unitary matrix in $\mathcal{M}_n(U\langle n \rangle)$.

Unless the quantum case, we do not suppose that $\bar{u} = (u_{ij}^*)_{1 \leq i, j \leq n}$ is unitary!

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Let $(U_t^{(d)})_{t \geq 0}$ be a Brownian motion on the (classical) unitary group $U(nd)$ (with n fixed). Then, we define the following process:

$$j_t^{(d)} : U\langle n \rangle \rightarrow L^{\infty-} \otimes \mathcal{M}_d$$

$$u_{ij} \mapsto [U_t^{(d)}]_{ij}$$

where $[.]_{ij}$ designates the (i, j) block of size $d \times d$ of a matrix.

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

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Where $L^{\infty-}$ are the random variable having moments of all order.

We say that $j_t^{(d)}$ is a quantum random variable on $U\langle n \rangle$.

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$ A stochastic process on
 $U\langle n \rangle$ Haar states on
Dual GroupsHaar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free caseIdea of the method for the
tensor case

Going further

Theorem (U, 2014; second proof Cebron, U, 2015)

When d goes to infinity, the process $(j_t^{(d)})_{t \geq 0}$ tends towards a free Lévy process $(j_t)_{t \geq 0}$ on $U\langle n \rangle$.

Remark

When $n = 1$, this is a classical result obtained by Biane in 1997.

Theorem (U, 2014)

The limit free Lévy process verifies

$$j_t : u_{ij} \mapsto \Psi_{ij}$$

where Ψ_t is a $n \times n$ matrix whose elements verify the free stochastic differential equation

$$d\Psi = \frac{i}{\sqrt{n}} dX_t \Psi_t - \frac{1}{2} \Psi_t dt$$

with X_t being a self-conjugate matrix whose elements are free additive Brownian motions (free equivalent of the GUE).

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U(n)$ A stochastic process on $U(n)$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Theorem

$$j_t : u_{ij} \mapsto \Psi_{ij}$$
$$d\Psi = \frac{i}{\sqrt{n}} dX_t \Psi_t - \frac{1}{2} \Psi_t dt$$

Remark

The form of the equation is very similar to the (classical) stochastic equation for the Brownian motion on $U(n)$:

$$dU_t = i dH_t U_t - \frac{1}{2} U_t dt$$

where H_t is the hermitian Brownian motion.

Theorem

$$j_t : u_{ij} \mapsto \psi_{ij}$$
$$d\psi = \frac{i}{\sqrt{n}} dX_t \psi_t - \frac{1}{2} \psi_t dt$$

Remark

The form of the equation is very similar to the (classical) stochastic equation for the Brownian motion on $U(n)$:

$$dU_t = i dH_t U_t - \frac{1}{2} U_t dt$$

where H_t is the hermitian Brownian motion.

Theorem (U, 2014)

The process $(j_t)_{t \geq 0}$ verifies a property called gaussianity (in the sense of Schürmann).

Remark

It is thus a good candidate to be called a Brownian motion

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Definition (Cebron, U, 2015)

We define a Haar state on $U\langle n \rangle$ as being a state h such that for any state ϕ :

$$(h \odot \phi)\Delta = h = (\phi \odot h)\Delta$$

with \odot being one of the five notions of independances (free, tensor, monotone, antimonotone, boolean).

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$ A stochastic process on
 $U\langle n \rangle$ Haar states on
Dual GroupsHaar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free caseIdea of the method for the
tensor case

Going further

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Remark

We need to restrict it to states because the independances are not necessarily distributive.

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

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Remark

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Remark

The boolean, monotone and antimonotone independances are defined on non-unital algebras. For our definition to make sense, we must consider $U\langle n \rangle = \mathbb{C} \oplus \text{Ker}\delta$.

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$ A stochastic process on
 $U(n)$ Haar states on
Dual GroupsHaar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free caseIdea of the method for the
tensor case

Going further

Theorem (Cebron, U, 2015)

For $n \geq 2$, there exists no tensor, free, boolean, monotone or anti-monotone Haar state.

(The obvious Haar weight on the unit circle works for $n = 1$ for all five independances)

Theorem (Cebron, U, 2015)

For $n \geq 2$, there exists no tensor, free, boolean, monotone or anti-monotone Haar state.

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Proof.

The ingredient of the proof is to build specific states ϕ and to find a contradiction... □

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Definition (Cebron, U, 2015)

We call Haar trace a tracial state that is absorbing for tracial states ϕ .

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Theorem (Cebron, U, 2015)

For $n \geq 2$, there exists no Haar trace for the boolean, monotone and anti-monotones independances, but there exists one for the free one and one for the tensor one.

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

We take (A, ψ) a noncommutative $*$ -probability space. We define:

$$\tilde{A} = E_{11}[A \sqcup M_n(\mathbb{C})]E_{11}$$

and

$$h = n[\psi * \text{tr}]$$

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U(n)$

A stochastic process on $U(n)$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

We take (A, ψ) a noncommutative $*$ -probability space. We define:

$$\tilde{A} = E_{11}[A \sqcup M_n(\mathbb{C})]E_{11}$$

and

$$h = n[\psi * \text{tr}]$$

We also define the map

$$j: \begin{array}{l} U\langle n \rangle \rightarrow \tilde{A} \\ u_{ij} \mapsto E_{1i} U E_{j1} \end{array}$$

with $U \in A$ a Haar unitary.

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Theorem (Cebron, U, 2015)

The state $h \circ j$ is a free Haar trace on $U\langle n \rangle$.

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$ A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

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Theorem (Cebron, U, 2015)

The state $h \circ j$ is a free Haar trace on $U\langle n \rangle$.

Proof.

Ingredients: Theorem of Nica and Speicher about compressed algebra
cumulants of Haar unitary
considering partitions.



Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

We consider the Hilbert space $H = l^2(\mathbb{Z}) \otimes \bigotimes_{k \in \mathbb{Z}} \mathcal{M}_n(\mathbb{C})$,
where the number of matrices different of I_n is finite.

We define the operator:

$$\begin{aligned} & U_{ij}(\delta_k \otimes (\dots \otimes M_{k-1} \otimes M_k \otimes M_{k+1} \otimes \dots)) \\ = & \delta_{k+1} \otimes (\dots \otimes M_{k-1} \otimes E_{ji} M_k \otimes M_{k+1} \otimes \dots) \\ & U_{ij}^*(\delta_k \otimes (\dots \otimes M_{k-1} \otimes M_k \otimes M_{k+1} \otimes \dots)) \\ = & \delta_{k-1} \otimes (\dots \otimes E_{ij} M_{k-1} \otimes M_k \otimes M_{k+1} \otimes \dots) \end{aligned}$$

And we set:

$$\Omega = \delta_0 \otimes \bigotimes_{k \in \mathbb{Z}} I_n$$

and $j : u_{ij} \mapsto U_{ij}$.

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$

A stochastic process on
 $U(n)$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Theorem (Cebron, U, 2015)

The state $h(a) = \langle \Omega, j(a)\Omega \rangle$ is the tensor Haar trace.

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$

A stochastic process on
 $U(n)$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

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Theorem (Cebron, U, 2015)

The state $h(a) = \langle \Omega, j(a)\Omega \rangle$ is the tensor Haar trace.

Proof.

It is easy to compute the values of h on words alternating in $*$. To compute more general words, we use

$$S_k = \left\{ l \in \{1, \dots, r\} : k = \#\{m > l : \epsilon_m = \emptyset\} - \#\{m \geq l : \epsilon_m = *\} \right\}$$

and h then factorizes according to these S_k . □

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$

A stochastic process on
 $U(n)$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Theorem (Cebron, U, 2015)

The free Haar trace is faithful, whereas the tensor one is not.

Introduction to the basic objects

Dual groups

The unitary Dual Group

A Brownian motion on $U\langle n \rangle$

A stochastic process on $U\langle n \rangle$

Haar states on Dual Groups

Haar states and independances

A negative answer

On the way towards a positive result

Haar trace

Idea of the method for the free case

Idea of the method for the tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U\langle n \rangle$

A stochastic process on
 $U\langle n \rangle$

Haar states on
Dual Groups

Haar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free case

Idea of the method for the
tensor case

Going further

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$ A stochastic process on
 $U(n)$ Haar states on
Dual GroupsHaar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free caseIdea of the method for the
tensor case

Going further

results already obtained:

- ▶ Replacing U by a free Lévy process U_t in our method, we get unitary quantum Lévy processes. We can write their Schürmann triple.

results already obtained:

- ▶ Replacing U by a free Lévy process U_t in our method, we get unitary quantum Lévy processes. We can write their Schürmann triple.

Question in process:

- ▶ What does it mean that we get only a Haar traces?
- ▶ What happens with other dual groups?

- ▶ Construction of a free Lévy process as high-dimensional limit of a Brownian motion on the unitary group, M.Ulrich, *Infin. Dim. Anal. Quant. Prob.*, 2015, vol. 18
- ▶ Haar states and Lévy processes on the unitary dual group, G.Cébron, M.Ulrich, *ArXiv:1505.08083*, to appear in *JFA*

Introduction to the
basic objects

Dual groups

The unitary Dual Group

A Brownian
motion on $U(n)$ A stochastic process on
 $U(n)$ Haar states on
Dual GroupsHaar states and
independances

A negative answer

On the way
towards a positive
result

Haar trace

Idea of the method for the
free caseIdea of the method for the
tensor case

Going further